

Fast Identification of HCF Properties by Self-Heating Measurements under Cyclic Loadings

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Abstract It is proposed to develop and identify a probabilistic two-scale model for HCF that accounts for the failure of samples but also for the thermal effects during cyclic loadings in a unified framework. The probabilistic model is based on a Poisson point process. Within the weakest link theory, the model corresponds to a Weibull law for the fatigue limits. The thermal effects can be described if one considers the same hypotheses apart from the weakest link assumption. A method of identification is proposed and uses temperature measurements to identify both mean fatigue limit and the scatter in an S/N curve. The validation of the model is obtained by predicting S/N curves for different metallic materials.

1 INTRODUCTION

The industry is one of the domains that needs good design procedures to ensure the reliability of structures. Since fracture of structures is often due to fatigue, probabilistic approaches have been developed in High Cycle Fatigue (HCF) life prediction. With these approaches, the scatter of fatigue data can be described, as well as the dependence of fatigue limits with the loaded volume and the type of loading by introducing the so-called effective volume. Consequently, the fatigue limits depend on material properties as well as the effective volume. The identification of probabilistic models is often based on numerous experimental fatigue tests to properly evaluate the scatter. The predictive capability of probabilistic models can be improved by identifying the origin of the scatter. In the present paper, the scatter of experimental fatigue results is assumed to be induced by a heterogeneous microplastic activity.

Some authors have worked on an estimation of the mean fatigue limit based upon temperature measurements [1-6]. This method consists in observing the thermal effect during cyclic loadings. It has the advantage of being performed in shorter time compared with traditional HCF tests. A good correlation can be obtained between the inception of a temperature increase and the mean fatigue limit. If a probabilistic approach useful for HCF life prediction is considered, then the analysis of the temperature measurements cannot be used in a classical way and needs to be revisited. It will be used to identify parameters relevant to a probabilistic model to describe the initiation stage in HCF.

The present paper is concerned with the derivation and identification of a probabilistic two-scale model for HCF that accounts for the failure of samples but also for the thermal effects during cyclic loadings in a unified framework. A two-scale model is used. It is assumed that HCF damage occurs at the microscopic scale and is caused by microplasticity that has a weak effect on the behavior of the material at the macroscopic scale. The development of the model is presented in two successive stages. In the first one, a deterministic viewpoint is chosen to show that the thermal effects can be accounted for by integrating the heat equation in which the microplastic dissipation is included. In view of a single microplastic hardening law, the empirical method proposed to estimate quickly the mean fatigue limits can be described to determine a deterministic (i.e., mean) fatigue limit. The model is then extended by introducing a random yield stress distribution at the microscopic scale. The new model describes more closely the thermal effects than the first one. An identification method of the scatter is proposed. It is based on the analysis of the thermal effects for the scatter and is applied on different metallic materials.

2 PRINCIPLE OF SELF-HEATING TESTS UNDER CYCLIC LOADING

The material studied herein is a ferrite-martensite dual-phase steel grade (around 15wt% of martensite, Arcelor designation: DP600) in its cold-rolled state. To evaluate the fatigue properties of the material, self-heating tests under cyclic loading can be used. They consist in applying successive series of 3000 cycles (Figure 1a) for different increasing stress amplitudes Σ_0 (Figure 1b). For each stress amplitude, the change of the temperature variation $\theta = T - T_0$ (where T is the current temperature of the sample during the test measured by a

thermocouple and T_0 the initial value) is recorded (Figure 1c). In the following developments, it is assumed that the temperature field is uniform in the sample. The mean temperature becomes stable after about 1000 cycles and equals $\bar{\theta}(\Sigma_0)$. Figure 1d shows the change of the steady-state mean temperature $\bar{\theta}(\Sigma_0)$ with the stress amplitude for two different tests. One can see that the information is reproducible. Furthermore, $\bar{\theta}(\Sigma_0)$ suddenly increases as the stress level is close to the fatigue limit of the material. This change of regime can be explained by the occurrence of the first slip bands on the surface [7]. In the next sections, two different models are introduced to relate the change of the mean temperature during series of cyclic loadings to the fatigue properties of the material.

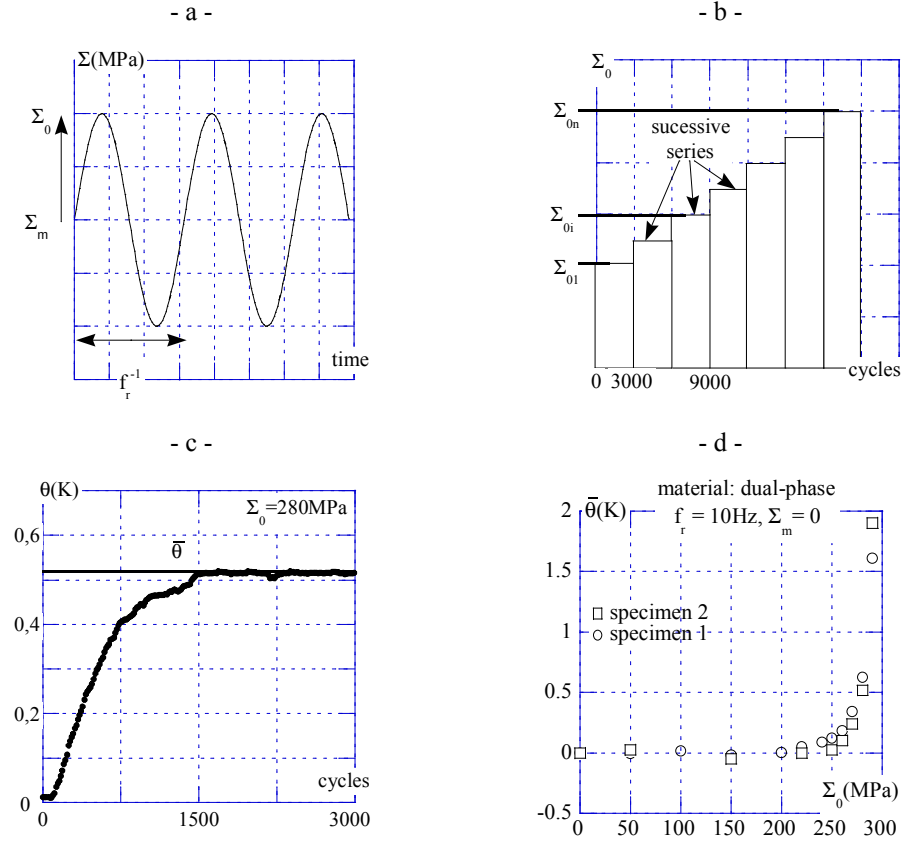


Figure 1. Self-heating test on a dual-phase DP600 steel: -a-cyclic loadings, -b-successive series of 3000 cyclic loadings for different increasing stress amplitudes Σ_0 , -c-change of the temperature during 3000 cyclic loadings for the same stress amplitude -d-change of the steady-state mean temperature with the load amplitude.

3 A DETERMINISTIC TWO-SCALE MODEL

The two-scale model used herein has been initially developed by Lemaitre and Doghri [8], and then extended by Lemaitre et al. [9] in a deterministic way. It is based on the hypothesis that HCF damage is localized at the microscopic scale, whose dimensions are smaller than those of the representative volume element (RVE) of volume V_{RVE} associated to the mesoscopic scale. Consequently, two phases are considered in the model. The first one is an inclusion of volume V_0 accounting for a set of grains where local plasticity occurs; the second one is the surrounding matrix describing the other grains. The two phases are assumed to have the same elastic behavior and the yield stress of the elasto-plastic inclusion σ_y^μ is less than that of the surrounding matrix that remains elastic. Let $\underline{\underline{\Sigma}}$ and $\underline{\underline{E}}$ be the mesoscopic stress and strain tensors, respectively, $\underline{\underline{\sigma}}$ and $\underline{\underline{\varepsilon}}$ the microscopic stress and strain tensors in the inclusions, respectively. An additive decomposition of strain with an elastic part $\underline{\underline{\varepsilon}}^e$ and a plastic part $\underline{\underline{\varepsilon}}^p$ is assumed for the inclusion

$$\underline{\underline{\varepsilon}} = \underline{\underline{\varepsilon}}^e + \underline{\underline{\varepsilon}}^p \quad (1)$$

A linear kinematic hardening is considered in this work. It follows that the back stress $\underline{\underline{X}}$ can be related to the plastic strain tensor $\underline{\underline{\epsilon}}^p$ in the inclusion by

$$\dot{\underline{\underline{X}}} = \frac{2}{3} C \dot{\underline{\underline{\epsilon}}}^p \quad (2)$$

where C is a parameter depending on the material, and a dotted variable corresponds to its first derivative with respect to time. A normality rule is assumed for the plasticity model

$$\dot{\underline{\underline{\epsilon}}}^p = \dot{\lambda} \frac{\partial f}{\partial \underline{\underline{S}}} \quad (3)$$

with the usual Kuhn-Tucker conditions, $\underline{\underline{S}} = \underline{\underline{\sigma}} - \frac{1}{3} \text{trace}(\underline{\underline{\sigma}}) \underline{\underline{I}}$ is the deviatoric stress tensor, $\underline{\underline{I}}$ the unit second order tensor, $\dot{\lambda}$ the plastic multiplier and f the yield surface

$$f = J_2(\underline{\underline{S}} - \underline{\underline{X}}) - \sigma_y^\mu \leq 0 \quad (4)$$

where J_2 is the second stress invariant, i.e., $J_2(\underline{\underline{S}} - \underline{\underline{X}}) = \sqrt{\frac{3}{2}(\underline{\underline{S}} - \underline{\underline{X}}) : (\underline{\underline{S}} - \underline{\underline{X}})}$, and ‘:’ the tensorial product contracted with respect to two indices. The microscopic stresses are evaluated from the mesoscopic stresses by means of a law of localization. This law is obtained by a homogenization procedure

$$\underline{\underline{\sigma}} = \underline{\underline{\Sigma}} - 2\mu(1 - \beta)(\underline{\underline{\epsilon}}^p) \quad (5)$$

with $\beta = \frac{2(4 - 5\nu)}{15(1 - \nu)}$ given by the Eshelby analysis of a spherical inclusion [10], μ and ν , the shear modulus and the Poisson’s ratio of the dual-phase steel, respectively.

The concepts of continuum thermodynamics are used to analyze the temperature measurements. At each instant t , the thermodynamic state of each RVE is characterized by state variables. In the present case, three state variables are used, namely, T the absolute temperature (or $\theta = T - T_0$), $\underline{\underline{\epsilon}}^e$ and $\underline{\underline{\epsilon}}^p$. The specific Helmholtz free energy of the RVE Ψ is defined by the following relationship

$$\Psi = \Psi_{\text{mat}}(1 - f_v) + \Psi_{\text{in}} f_v \quad (6)$$

where $f_v = V_0 / V_{\text{RVE}}$ is the volume fraction of the inclusion in the RVE, Ψ_{mat} and Ψ_{in} are the free energy densities of the matrix and the inclusion, respectively

$$\rho \Psi_{\text{mat}} = \frac{1}{2} \underline{\underline{E}} : \underline{\underline{K}} : \underline{\underline{E}} - \alpha \theta \underline{\underline{I}} : \underline{\underline{K}} : \underline{\underline{E}} - \frac{c \theta^2}{2T_0} \quad (7)$$

$$\rho \Psi_{\text{in}} = \frac{1}{2} \underline{\underline{\epsilon}}^e : \underline{\underline{K}} : \underline{\underline{\epsilon}}^e + \frac{1}{2} \underline{\underline{\epsilon}}^p : \left(\frac{2}{3} C \underline{\underline{I}} \right) : \underline{\underline{\epsilon}}^p - \alpha \theta \underline{\underline{I}} : \underline{\underline{K}} : \underline{\underline{\epsilon}}^e - \frac{c \theta^2}{2T_0} \quad (8)$$

where $\underline{\underline{K}}$ is the elastic tensor, ρ the mass density, c the specific heat capacity, $\underline{\underline{I}}$ the unit fourth order tensor, and α the coefficient of thermal expansion. In Eqn. (7), the free energy characterizes a thermo-elastic behavior for which the first term is the elastic or recoverable contribution, the second corresponds to the thermo-elastic coupling and the third to the temperature contribution associated to temperature variations that remain small (with $\theta < 10\text{K}$). In the right-hand side of Eqn. (8), in addition to the previous contributions, the second term corresponds to the stored energy caused by the hardening of the material at the scale of the inclusion.

By application of the first law of thermodynamics, the external power Π can be written as

$$\Pi = (1 - f_v) \underline{\underline{\Sigma}} : \dot{\underline{\underline{E}}} + f_v \underline{\underline{\sigma}} : \dot{\underline{\underline{\epsilon}}}^e + f_v (\underline{\underline{\sigma}} - \underline{\underline{X}}) : \dot{\underline{\underline{\epsilon}}}^p + f_v \underline{\underline{X}} : \dot{\underline{\underline{\epsilon}}}^p \quad (9)$$

where $(1-f_v)\underline{\underline{\Sigma}}:\underline{\underline{\dot{E}}}$ and $f_v\underline{\underline{\sigma}}:\underline{\underline{\dot{\epsilon}^p}}$ are the recoverable elastic power, $f_v\underline{\underline{X}}:\underline{\underline{\dot{\epsilon}^p}}$ the plastic power and $f_v(\underline{\underline{\sigma}}-\underline{\underline{X}}):\underline{\underline{\dot{\epsilon}^p}}$ the dissipated power. The dissipated power Δ describes, among other things, the losses in heat during the motion of dislocations and is given by

$$\Delta = f_v(\underline{\underline{\sigma}} - \underline{\underline{X}}):\underline{\underline{\dot{\epsilon}^p}} = f_v(\underline{\underline{\sigma}} - \underline{\underline{X}}):\dot{\lambda}\frac{\partial f}{\partial \underline{\underline{\sigma}}} \quad (10)$$

The plastic multiplier $\dot{\lambda}$ is obtained from the consistency condition

$$\dot{f} = \frac{\partial f}{\partial \underline{\underline{\sigma}}}:\underline{\underline{\dot{\sigma}}} + \frac{\partial f}{\partial \underline{\underline{X}}}:\underline{\underline{\dot{X}}} = 0 \quad (11)$$

By using the law of localization (5) and the hardening kinetics (2), Eqn. (11) can be rewritten as

$$\frac{\partial f}{\partial \underline{\underline{\sigma}}}:\left[\underline{\underline{\dot{\Sigma}}} - 2\mu(1-\beta)\underline{\underline{\dot{\epsilon}^p}} - \frac{2}{3}C\underline{\underline{\dot{\epsilon}^p}}\right] = 0 \quad (12)$$

Consequently, one can deduce the following relationship by using the normality rule (3) with the hardening modulus $h = C + 3\mu(1-\beta)$

$$\dot{\lambda} = \frac{H(f)}{h} \frac{\partial f}{\partial \underline{\underline{\sigma}}}:\underline{\underline{\dot{\Sigma}}} \quad (13)$$

where H is the Heaviside step function (i.e., $H(f)=0$ if $f < 0$ and $H(f)=1$ if $f = 0$). The dissipated energy density D for a loading cycle becomes

$$D = \int_{\text{cycle}} f_v \frac{H(f)}{h} \left(\frac{\partial f}{\partial \underline{\underline{\sigma}}}:\underline{\underline{\dot{\Sigma}}} \right) \left(\underline{\underline{\sigma}} - \underline{\underline{X}} \right) \frac{\partial f}{\partial \underline{\underline{\sigma}}} dt \quad (14)$$

The integration of this last equation yields

$$D = \frac{4f_v\sigma_y^\mu}{h} \langle \underline{\underline{\Sigma}}_0 - \underline{\underline{\sigma}}_y^\mu \rangle \quad (15)$$

where $\langle . \rangle$ are the Macauley brackets (i.e. positive part of '.'). The stored power Π_s is expressed as

$$\Pi_s = f_v \underline{\underline{X}}:\underline{\underline{\dot{\epsilon}^p}} = f_v \frac{2}{3} C \underline{\underline{X}}:\underline{\underline{\dot{X}}} \quad (16)$$

so that the stored energy for one cycle vanishes. Consequently, the only energy that accumulates during the cycles is the dissipated part. An energetic initiation criterion is chosen by assuming for the material the existence of a critical dissipated energy (E_c) above which failure occurs. For each cycle, the dissipated energy is constant. Consequently, the number of cycles to failure N is related to D and E_c by

$$ND = E_c \quad (17)$$

so that the number of cycles to failure can be linked to the stress level $\underline{\underline{\Sigma}}_0$ and the microscopic yield stress $\underline{\underline{\sigma}}_y^\mu$ through Eqn. (15) by

$$N = \frac{A}{\langle \underline{\underline{\Sigma}}_0 - \underline{\underline{\sigma}}_y^\mu \rangle \underline{\underline{\sigma}}_y^\mu} \quad (18)$$

where $A = hE_c/4f_v$ is a parameter depending on the considered material. Let $\underline{\underline{\sigma}}_y^\mu = \bar{\underline{\underline{\Sigma}}}_\infty$ denotes the mean fatigue limit, Figure 2 shows that the model agrees well with experimental results by using a relevant couple of parameters ($A, \underline{\underline{\sigma}}_y^\mu$).

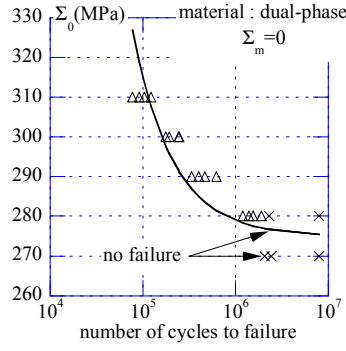


Figure 2. S/N curve of the studied dual-phase steel ($\Sigma_m = 0$).

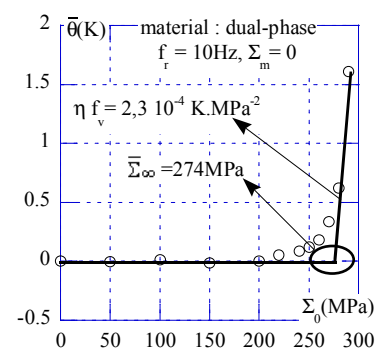


Figure 3: identification of η and $\bar{\Sigma}_\infty$ for the studied dual-phase steel ($f_r = 10$ Hz, $\Sigma_m = 0$).

The previous results are now used to analyze the heat transfer induced by the intrinsic dissipation D . The heat equation is deduced from the two first principles of thermodynamics. Let us assume that the convective terms of the material time derivative and the coupling between temperature and the hardening variable are negligible, and that the external heat supply r is time-independent. The local heat conduction equation is given by [11]

$$\rho c \dot{\theta} - k \Delta \theta = D + T \rho \frac{\partial^2 \Psi}{\partial T \partial \underline{\epsilon}} : \dot{\underline{\epsilon}}^e = \rho c S_t \quad (19)$$

where $\rho c \dot{\theta}$ denotes the absorption heat capacity and $k \Delta \theta = -\text{div}(\vec{q})$ denotes the losses by conduction, with k the thermal conductivity, $\vec{q}$ the heat influx vector and S_t the heat source depending on the thermoelastic coupling and the intrinsic dissipation. For small temperature variations ($\theta < 10$ K), k and c are independent of the temperature. If the temperature is uniform, an analysis of the thermal boundary conditions with the surrounding air and with the grips of the testing machine gives the following relationship for the mean losses by conduction [6]

$$\frac{1}{V} \int \text{div}(\vec{q}) dV = \left(\frac{2h_1(e+1)}{el} + \frac{2h_2}{L} \right) \theta = \frac{\rho c}{\tau_{eq}} \theta \quad (20)$$

where h_1 and h_2 are respectively the heat exchange parameters between the surrounding air and the specimen, and between the grips and the sample. The characteristic time τ_{eq} depends on the thermal boundary conditions. The local heat conduction equation becomes

$$\dot{\theta} + \frac{\theta}{\tau_{eq}} = S_t \quad (21)$$

Let us recall that the mean value per cycle of the temperature is sought. Consequently, the mean heat source has to be evaluated to solve the heat conduction equation. The value of the mean thermoelastic contribution vanishes. The mean heat source reads

$$\rho c S_t = f_r D \quad (22)$$

where f_r is the loading frequency, and the mean steady-state temperature $\bar{\theta}$ is the particular solution of the differential equation (21)

$$\bar{\theta} = \frac{4f_v f_r \tau_{eq} \sigma_y^\mu}{h \rho c} \langle \Sigma_0 - \sigma_y^\mu \rangle \quad (23)$$

Equation (23) shows that below $\bar{\Sigma}_\infty = \sigma_y^\mu$, no temperature variation is observed. Above the latter, a linear response of the temperature as a function of the stress is obtained whose slope $(4f_r f_v \tau_{eq} / h \rho c) \sigma_y^\mu = \eta f_v \sigma_y^\mu$ where $\eta = 4f_r \tau_{eq} / h \rho c$ is linearly related to the frequency f_r of the cycles. Figure 3 shows the result of the identification of the two parameters $\bar{\Sigma}_\infty$ and ηf_v of the model. A good agreement is obtained in accordance with previous results obtained on other materials [1, 4]. However, it can be noted that the transition defined by a

unique threshold $\bar{\Sigma}_\infty$ cannot be determined very accurately. An inaccuracy exists on the value the fatigue limit when these measurements are used to determine the mean fatigue limit. This is due to the fact that there is a gradual temperature increase instead of a two-line trend. Yet it allows for a first estimate of a fatigue limit in a very short time duration. As shown in Figure 2, there exists a scatter in fatigue data that needs to be accounted for. The following section deals with an extension of the present approach by considering that the microscopic yield stresses are scattered.

4 A PROBABILISTIC TWO-SCALE MODEL

Let us assume that the microscopic yield stress is a probabilistic variable. An active site (in terms of fatigue mechanism) is defined as an inclusion whose equivalent stress is greater than the microscopic yield stress. These sites are assumed to be described by a Poisson point process. This process is specified by an intensity function λ that is assumed to follow a power law of the equivalent stress amplitude Σ

$$\lambda(\Sigma) = \frac{1}{V_0} \left(\frac{\Sigma}{S_0} \right)^m \quad (24)$$

where $V_0 S_0^m$ and m are two parameters depending on the material. The probability P_k of finding k active sites in a domain Ω of volume V_Ω follows a Poisson distribution

$$P_k(V_\Omega) = \exp(-\lambda V_\Omega) \frac{(\lambda V_\Omega)^k}{k!} \quad (25)$$

and the average number of active sites $\bar{n}_a$ in a domain Ω of volume V_Ω can be related to the intensity of the Poisson process λ by

$$\bar{n}_a = \lambda(\Sigma) V_\Omega \quad (26)$$

It follows that the intensity of the Poisson process corresponds to the *mean* density of active sites for a given applied stress amplitude. This model can be used to describe HCF results. Within the weakest link theory, the failure probability P_F of a domain Ω of volume V_Ω under a homogeneous stress amplitude Σ is equal to the probability of finding at least one active inclusion in the volume V_Ω

$$P_F = P_{k \geq 1}(V_\Omega) = 1 - \exp(-\lambda V_\Omega) \quad (27)$$

By using the intensity λ defined in Eqn. (24), the failure probability reads

$$P_F = 1 - \exp \left[- \frac{V_\Omega}{V_0} \left(\frac{\Sigma}{S_0} \right)^m \right] \quad (28)$$

Equation (28) corresponds to a Weibull model [12, 13] that can be recovered within the framework of a Poisson point process coupled with a weakest link assumption. The parameters of the intensity λ (24) then correspond to the Weibull parameters. When the stress field is no longer homogeneous, the previous result can be generalized

$$P_F = 1 - \exp \left[- \frac{V_\Omega H_m}{V_0} \left(\frac{\Sigma_F}{S_0} \right)^m \right] \quad (29)$$

where $H_m = \frac{1}{V_\Omega} \int \left(\frac{\Sigma}{\Sigma_F} \right)^m dV$ is the stress heterogeneity factor and $\Sigma_F = \max_\Omega(\Sigma)$. The effective volume is the product of the loaded volume V_Ω by the stress heterogeneity factor H_m . The HCF life prediction is often characterized by the mean fatigue limit $\bar{\Sigma}_\infty$ and the corresponding variance $\bar{\Sigma}_\infty^2$ that are related to the Weibull parameters m and $V_0 \sigma_0^m$ by

$$\bar{\Sigma}_{\infty} = S_0 \left(\frac{V_0}{V_{\Omega} H_m} \right)^{\frac{1}{m}} \Gamma \left(1 + \frac{1}{m} \right) \quad (30)$$

$$\bar{\Sigma}_{\infty}^2 = S_0^2 \left(\frac{V_0}{V_{\Omega} H_m} \right)^{\frac{2}{m}} \left(\Gamma \left[1 + \frac{2}{m} \right] - \Gamma^2 \left[1 + \frac{1}{m} \right] \right) \quad (31)$$

where $\Gamma(p) = \int_0^{\infty} t^{p-1} \exp(-t) dt$ is the gamma function.

Equation (30) accounts for scale effects associated to the Weibull model. Namely, the higher the loaded volume, the smaller the mean fatigue limit; the higher the stress heterogeneity (i.e., the smaller H_m), the higher the mean fatigue limit. These coupled effects are described by the introduction of the so-called effective volume $V_{eff} = V_{\Omega} H_m$. Consequently corrections are to be performed when comparing data of samples tested in tension/compression to those of structures for which the stress field is different as well as the loaded volume. This analysis is relevant when initiation occurs within the volume as observed for dual-phase steels. Had initiation occurred at or close to the surface, an effective surface instead of an effective volume should have been used.

In the following analysis, the heat transfer analysis is performed again. However, the weakest link hypothesis is not considered since it is assumed that the gradual temperature change is associated with the gradual microplastic dissipation corresponding to a random distribution of yield stresses. The underlying Poisson point process as well as the power law of its intensity are still used. With a Poisson point process, $\frac{d\lambda}{d\Sigma} V_{\Omega} d\Sigma$ is the mean number of inclusions activated between an equivalent stress Σ and $\Sigma + d\Sigma$ in a domain of volume V_{Ω} , i.e. the number of sites whose mean fatigue limit lies between Σ and $\Sigma + d\Sigma$, and whose dissipated energy $V_{RVE} D(\Sigma)$ during a load cycle is given by

$$V_{RVE} D(\Sigma) = \frac{4V_0}{h} \Sigma \langle \Sigma_0 - \Sigma \rangle \quad (32)$$

For a specimen of volume V_{Ω} , the global dissipated energy density $\bar{\Delta}$ for an equivalent stress amplitude Σ_0 during a load cycle is expressed as

$$\bar{\Delta} = \int_0^{\Sigma_0} V_{RVE} D(\Sigma) \frac{d\lambda}{d\Sigma} d\Sigma = \frac{4V_0}{h} \frac{m}{(m+1)(m+2)} \frac{\Sigma_0^{m+2}}{(V_0^{1/m} S_0)^m} \quad (33)$$

The change of the mean steady-state temperature is obtained by the resolution of the heat conduction equation (21) by using $\bar{\Delta}$ instead of D

$$\bar{\theta} = \eta V_0 \frac{m}{(m+1)(m+2)} \frac{\Sigma_0^{m+2}}{(V_0^{1/m} S_0)^m} \quad (34)$$

Equation (34) does not show a sudden change of the mean steady-state temperature with the stress amplitude but rather a gradual increase in accordance with Fig. 5. The present model is dependent on three parameters, namely, ηV_0 , $S_0 (V_0)^{1/m}$ and m .

5 IDENTIFICATION PROCEDURE AND PREDICTIONS

In the following, the different parameters are identified: namely, $S_0 (V_0)^{1/m}$ characterizes the level of the mean fatigue limit for an effective volume equal to V_0 , ηV_0 is a parameter depending on material to describe the temperature kinetics, m measures the scatter of the fatigue limits, and A characterizes the initiation criterion. First, the parameters A and $\bar{\Sigma}_{\infty}$ are identified by using the classical Wöhler curve and Stromeier's law

$$N = \frac{A}{\langle \Sigma_0 - \bar{\Sigma}_{\infty} \rangle \bar{\Sigma}_{\infty}} \quad (35)$$

where 15 samples are usually needed. Then the values for m and ηV_0 are determined from the analysis of the temperature measurements (Eqn. (34)) for which an additional sample is used. The results of the identification are shown in Fig. 4. The temperature measurements are better described than with the deterministic approach (Fig. 5), especially in the transition regime. The scale parameter $S_0(V_0)^{1/m}$ is obtained by using the expression (30) of the mean fatigue limit for samples with an effective volume equal to $V_{eff} = 620 \text{ mm}^3$.

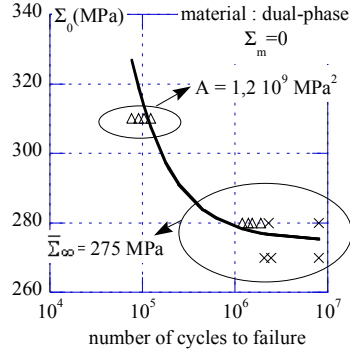


Figure 4: identification of A and $\bar{\Sigma}_\infty$ from the classical Wöhler curve.

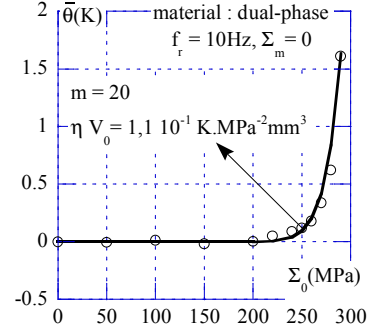


Figure 5: identification of ηV_0 and m from temperature measurements.

The validations concern the prediction of, on the one hand, the scatter of the experimental fatigue result in tension/compression (load ratio $R = -1$) and, on the other hand, the Wöhler curve for the same dual-phase steel loaded in alternate bending ($R = -1$). Figures 6 and 7 show the comparison between the experimental results obtained by classical fatigue tests and the prediction of the model for three different failure probabilities. The Wöhler curves can be predicted in a reasonable way.

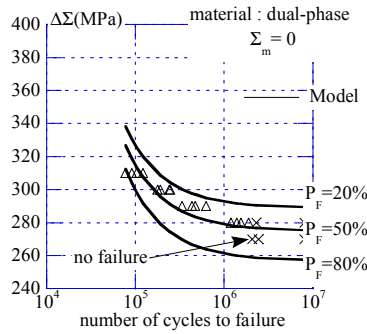


Figure 6: comparison between the prediction of the probabilistic model and the experimental results.

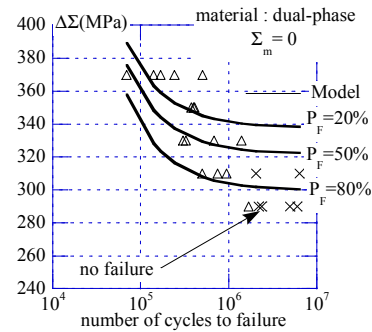


Figure 7: comparison between the prediction of the probabilistic model and the experimental results.

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